

An 80 MeV Injector Linac for BSF Future Project

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1. Introduction

A rapid cycle proton synchrotron with a maximum kinetic energy 800 MeV, maximum intensity 6×10^{13} ppp and repetition rate 50 Hz is under consideration as a BSF future project¹⁾. The injector will be 80 ~ 100 MeV H^- linac with pulse width 350 μs and intensity 30 mA.

The success²⁾ of RFQ in LANL suggests the alternative to a huge Cockcroft-Walton preinjector. Even though the attainable maximum energy through RFQ requires a deliberate study, 1 MeV will certainly be obtainable, which enables the use of klystron around 400 MHz as RF power source. Therefore the accelerating structure of this injector is divided into three stages: 50 ~ 100 keV H^- ion source, 1 MeV RFQ or APF and 80 MeV Alvarez linac. This paper describes the design study, mainly of Alvarez structure part.

2. Computation of Parameters by the Aid of SUPERFISH

If an Alvarez linac is excited at 400 MHz, the diameter of the cavity reduces nearly to the half of the present proton linac³⁾ at KEK and the average accelerating field can be made higher, say 3.5 MV/m. For a flat faced drift tube with fixed bore diameter and corner radius ($A = 1.0$ cm, $RHC = 0.5$ cm and $RH = 1$ cm, Fig. 1),

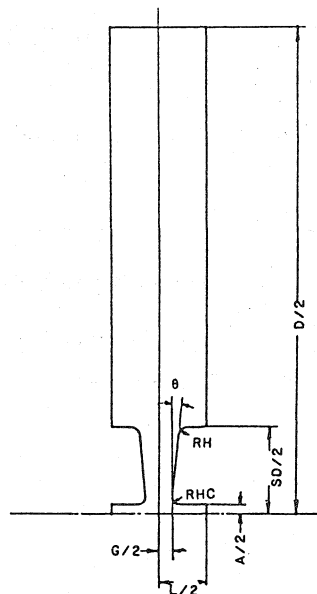


Fig. 1

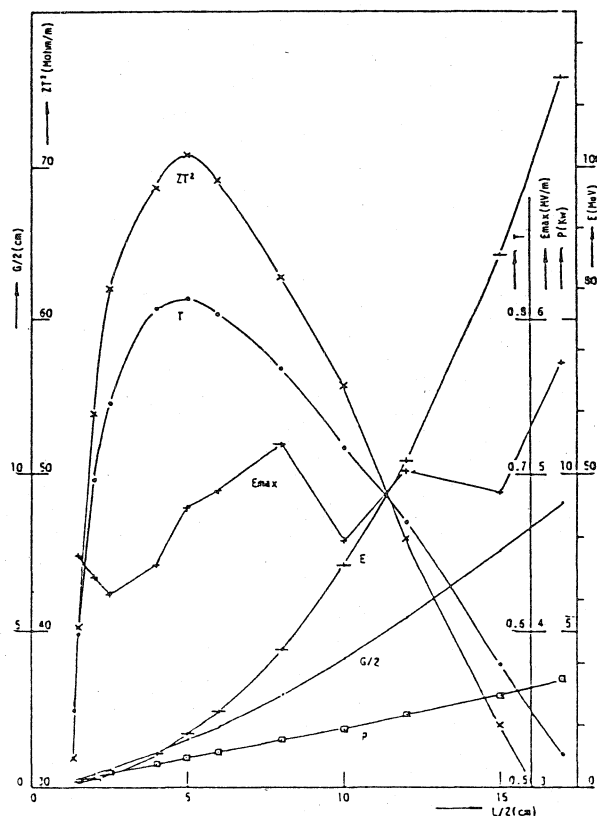


Fig. 2

an optimization was made at relatively low energy to obtain drift tube outer diameter $SD = 9.0$ cm and cavity diameter $D = 48$ cm. Next for these fixed dimensions, the half gap length ($L/2$) was computed that will give the resonant frequency of 400 MHz at different representative cell length. The values of transit time factor (T), effective shunt impedance (ZT^2), power loss (P), etc. were tabulated as the function of $L/2$ and used for interpolation (Fig. 2).

3. Cell Dimension and Longitudinal Oscillation

The geometrical and RF parameters for a constant synchronous phase ($\phi_s = -30^\circ$) are computed for several average accelerating field. The cell number, total length and RF power loss at $\bar{E}_0 = 3.0$ MV/m, 3.5 MV/m and 4.0 MV/m are shown in Table I. Table II shows the energy, cavity length, number of cells contained in a cavity and RF power for 5 partitions at $\bar{E}_0 = 3.5$ MV/m. From the view point of cavity numbers, cell numbers, RF power and sparking limit (the surface electric field is less than one third of Kilpatrick criterion), $\bar{E}_0 = 3.5$ MV/m is chosen for the first candidate.

\bar{E}_0 (MV/m)	Cell Number	Total length (m)	RF Power Loss (MW)
3.0	252	43.86	3.76
3.5	216	37.59	4.39
4.0	189	32.89	5.01

Table I

80 MeV Alvarez Linac (5 Cavities)

Cavity Number	Energy (MeV)	Length (m)	Number of Cells	Power (MW)		
				Cavity	Beam	Total
1	18.8602	7.4444	82	0.869	0.566	1.434
2	35.9399	7.5470	43	0.884	0.512	1.396
3	51.6727	7.5034	34	0.899	0.472	1.371
4	66.1500	7.3513	29	0.910	0.434	1.344
5	80.1218	7.5505	27	0.939	0.419	1.358

Table II

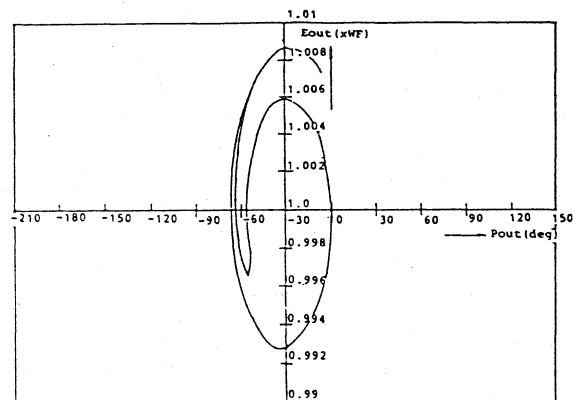
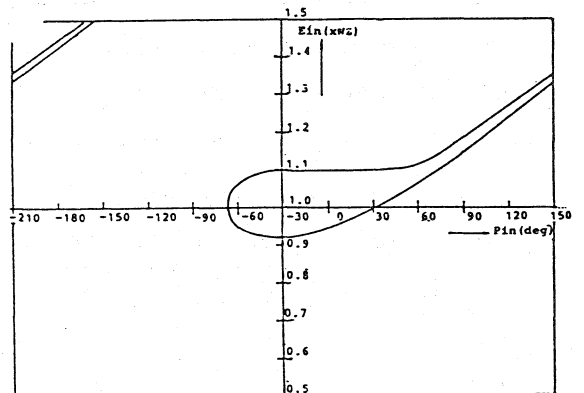


Figure 3(a) is the acceptance in ϕ - E space, in which WZ is the synchronous injection energy (1 MeV). As is well known a beam with several percent higher energy than WZ can be better captured. Figure 3(b) is the emittance, showing that the energy spread of the beam at the exit extends from -0.7% to $+0.6\%$ of the synchronous energy WF (80.1 MeV) and is

Fig. 3(a)(top) and (b)(bottom)

bunched within $\pm 30^\circ$. As this acceptance corresponds to the beam of Fig. 3(a), the actual one can be expected to be much smaller. Figure 4 shows the phase oscillation from 1 MeV to 3 MeV and particle distribution in the phase-energy space at the exit (upper right).

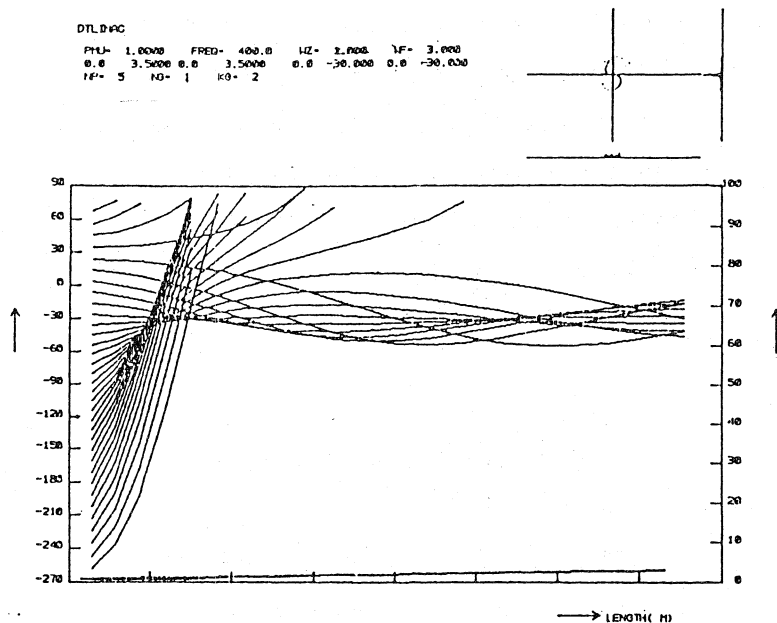


Fig. 4

4. Focussing System

As is the case of KEK 20 MeV linac⁴⁾, the field gradient $G(N)$ of the N -th quadrupole magnet is so designed as to make $\Lambda (= |K(S)| \cdot L_{\text{eff}}/L)$ constant. Therefore,

$$G(N) = (-1)^{N+1} G(0) \frac{\beta_{\text{in}} \cdot \gamma(N)}{\beta(N) \gamma_{\text{in}}} \left(\frac{L}{L_{\text{eff}} N} \right) \cdot \frac{1}{2}$$

in FDFD scheme, where L and L_{eff} means the cell length and effective quadrupole magnet length, respectively. β and γ are relativistic factors. The defocussing effects due to RF and space-charge are taken into consideration as impulse effects at the center of drift space. Figure 5 shows the normalized acceptance for the cell structure described in Section 3.

The abscissa designates the value of field gradient of the first quadrupole

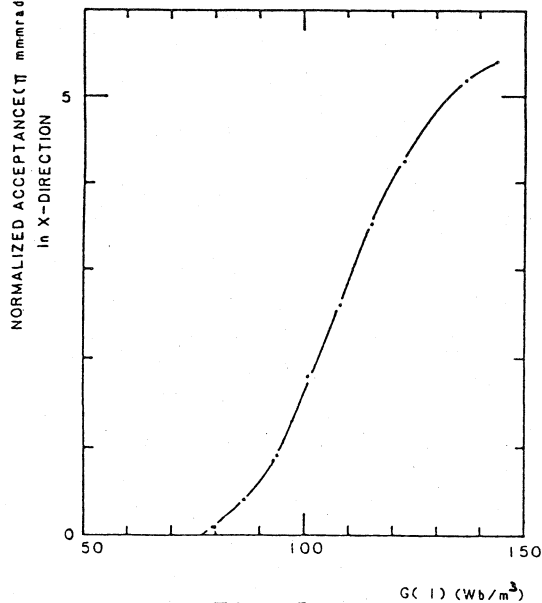


Fig. 5

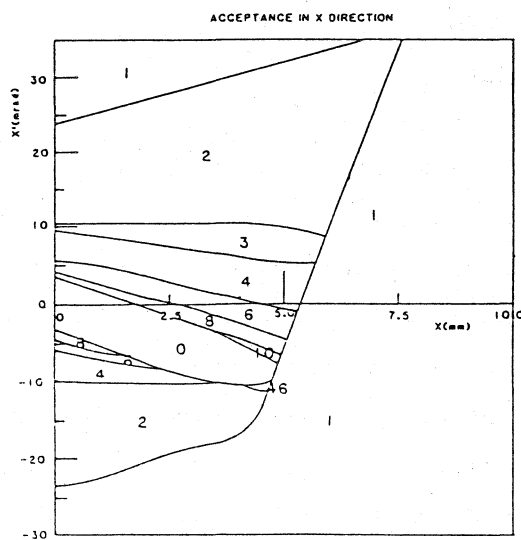


Fig. 6

magnet. Due to a relatively high accelerating field $\bar{E}_0 = 3.5 \text{ MV/m}$, the defocussing force becomes so large that focussing field gradient $G(1)$ is critical above 100 Wb/m^3 . This will be overcome by using rare earth-cobalt permanent magnet⁵). The acceptance at $G(1) = 130 \text{ Wb/m}^3$ is shown in Fig. 6, where the number indicates the cell where the beam is lost and the beam with number "0" is accepted throughout the structure.

5. Discussion

If the face of a drift tube is tapered ($\theta \neq 0$ in Fig. 1), the transit time factor and effective shunt impedance increases remarkably (Fig. 7). The electric field on the surface is shown in Fig. 8. These results suggest to use tapered drift tubes with 15° for $15 \sim 50 \text{ MeV}$, 30° for $50 \sim 80 \text{ MeV}$ and 45° for $80 \sim 100 \text{ MeV}$ as an illustration.

The build up time of the cavity for $Q = 6.5 \times 10^4$ is about $275 \mu\text{s}$. If one adds $100 \mu\text{s}$ as a margin to the minimum beam duration $350 \mu\text{s}$, the duty factor becomes 3.6 %. 400 MHz klystron with duty factor 5 % and peak power 2.5 MW is expected to be developed without much difficulty. Figure 9 shows the block diagram of the first scheme described in Section 3. Using tapered faced drift tubes, however, 100 MeV may be obtained by the same number of klystrons.

The experimental studies of post couplers and permanent

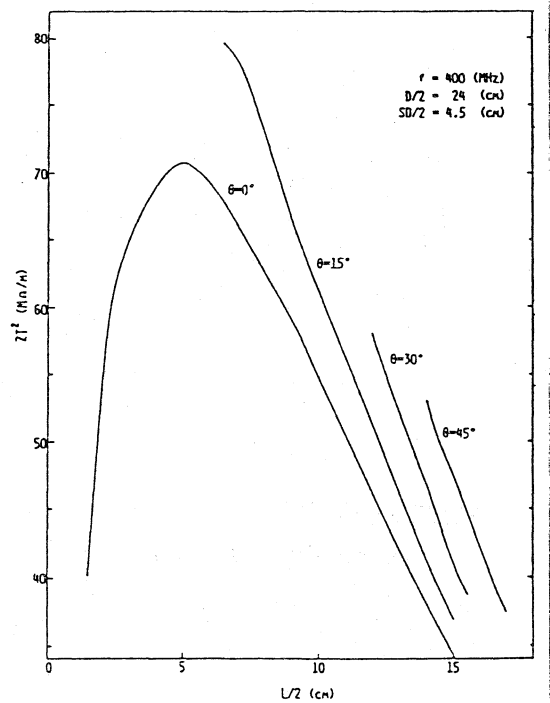


Fig. 7

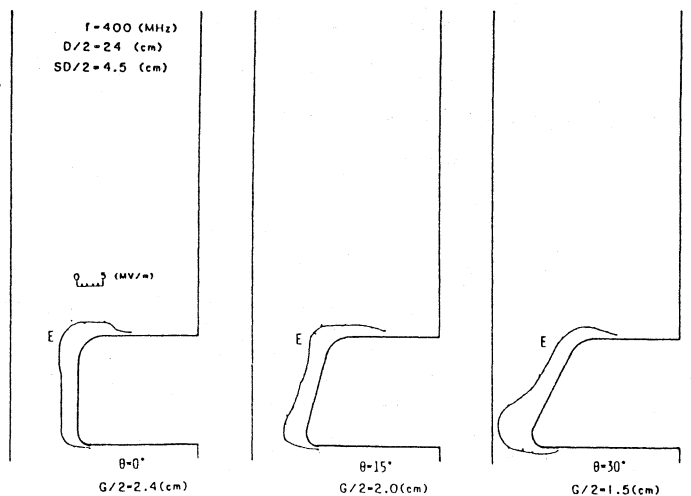
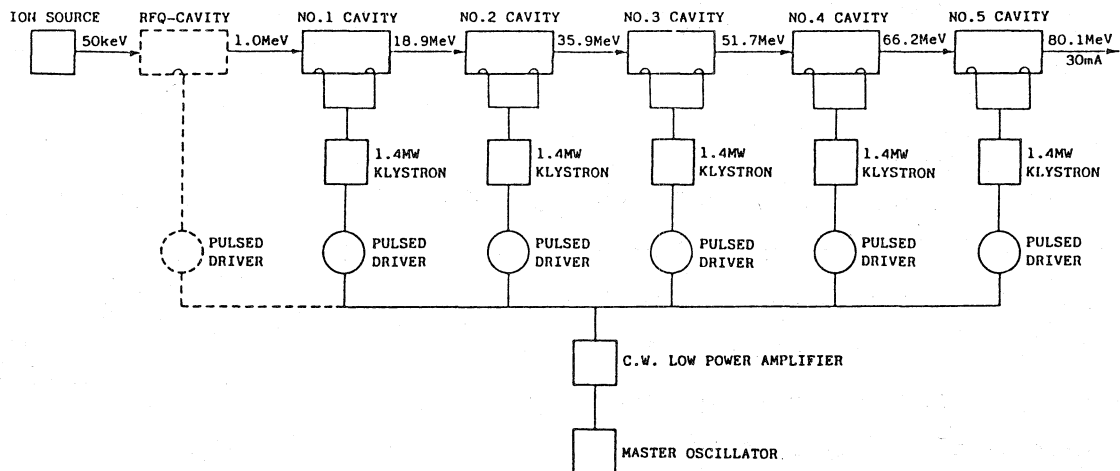


Fig. 8

magnets are essential to realize the present proposals. The investigations of low energy structure, RFQ or APF, and H^- ion sources need to be started later on. The intensity modulation of the beam with the frequency of synchrotron at an early stage is proposed to avoid the radioactive contamination of the synchrotron.



OUTLINE OF 80MeV INJECTOR LINAC

Fig. 9

Acknowledgement

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References

1. Proceedings of the Meeting on BSF Future Prospects, KEK Internal 82-6 (1982) (in Japanese).
2. R. H. Stokes, T. P. Wangler and K. R. Crandall, "The Radio-frequency Quadrupole - A New Linear Accelerator", IEEE Trans. Nucl. Sci. NS-28, 1999 (1981).
3. H. Matsumoto, J. Tanaka, I. Sato, S. Okumura and S. Inagaki, "Synchronous Phase Law Experiment in the KEK Linac", KEK-77-3 (1977).
4. M. Kobayashi, "Design of the Focussing System for KEK Linear Accelerator", KEK-73-4 (1973).
5. T. Takenaka and S. Inagaki, "Test Fabrication of a Samarium-cobalt Quadrupole Magnet", Proc. 5th Linear Accelerator Meeting in Japan, 169 (1980) (in Japanese).