

MOOT06

レーザー駆動誘電体加速器の設計

DESIGNING OF THE LASER DRIVEN DIELECTRIC ACCELERATOR

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Low-dose Radiation Risk

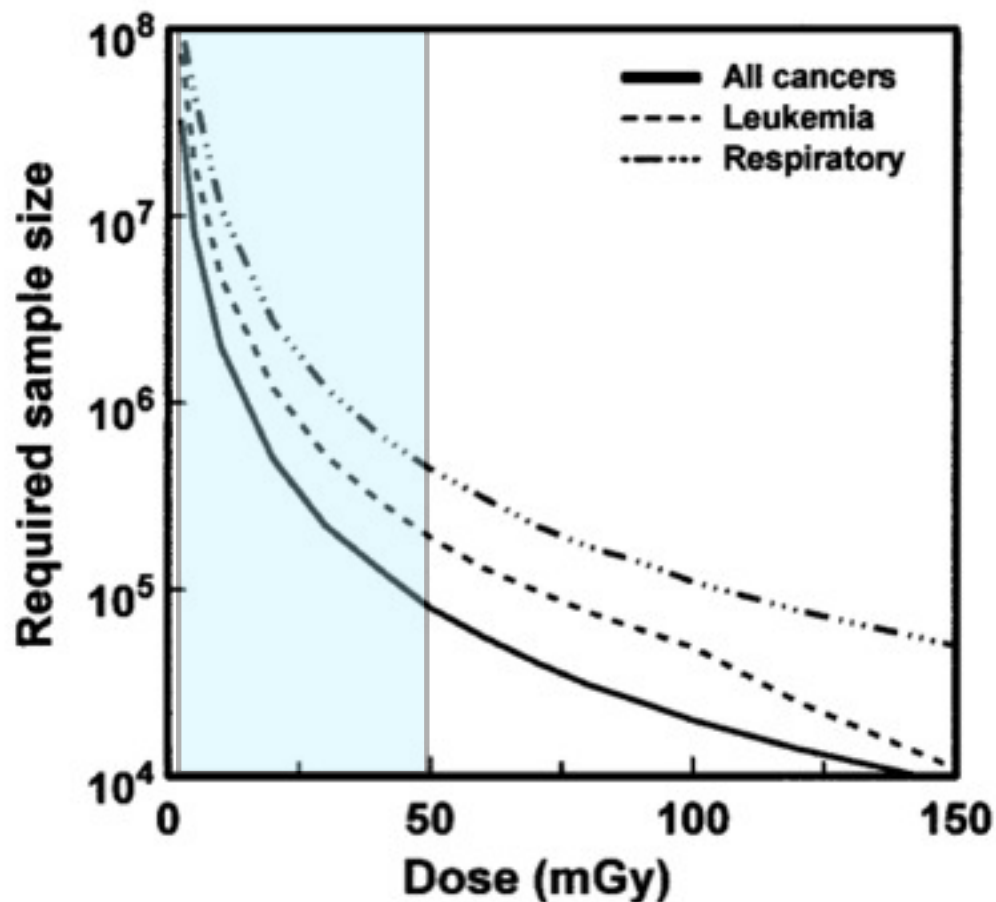


Fig. 1. Size of a cohort exposed to different radiation doses, which would be required to detect a significant increase in cancer mortality in that cohort, assuming lifetime follow-up (9).

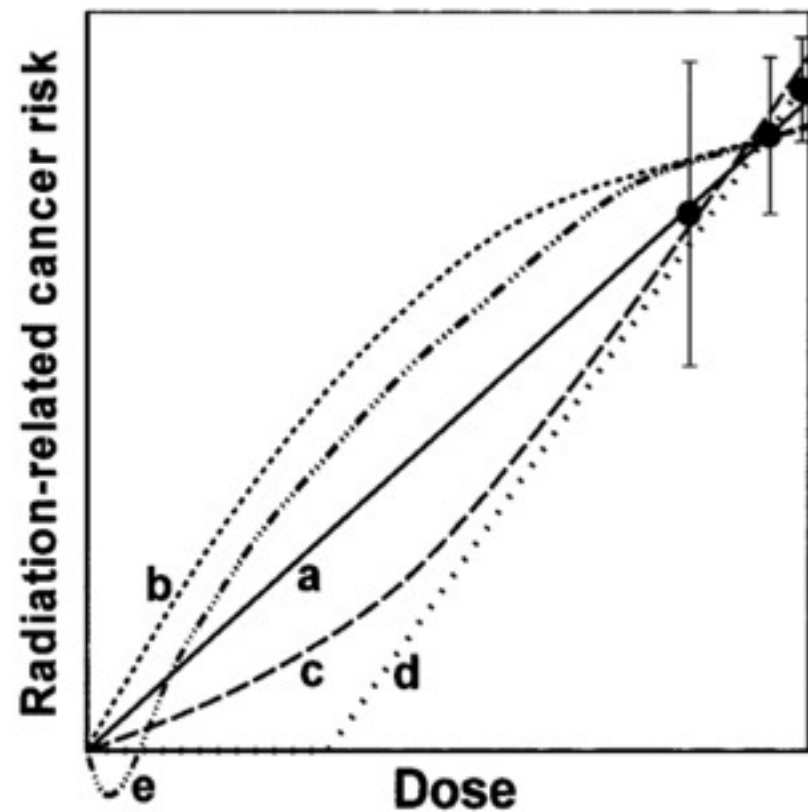


Fig. 3. Schematic representation of different possible extrapolations of measured radiation risks down to very low doses, all of which can in principle, be consistent with higher-dose epidemiological data. Curve a, linear extrapolation; curve b, downwardly curving (decreasing slope); curve c, upwardly curving (increasing slope); curve d, threshold; curve e, hormetic.

Brenner, D. J., et al., "Cancer risks attributable to low doses of ionizing radiation: Assessing what we really know", PNAS **100**, 13765 (2003).

Requirements

Electron / X-ray is much easier to make the compact source.

Energy deposition ; LET(linear energy transfer) e/x < ion
LET(keV/ μm) <1 ; >100

The beam size (irradiation area) is smaller than 1 μm .

The charge is 0.01 fC to 0.1 fC. ($\approx 10^2 - 10^3$ electrons/ μm^2)

The beam energy is several tens keV to 1 MeV.

The pulse width is sub fs. (μs -bunch)

The exit of the accelerator and specimen are observed through the transparent window.

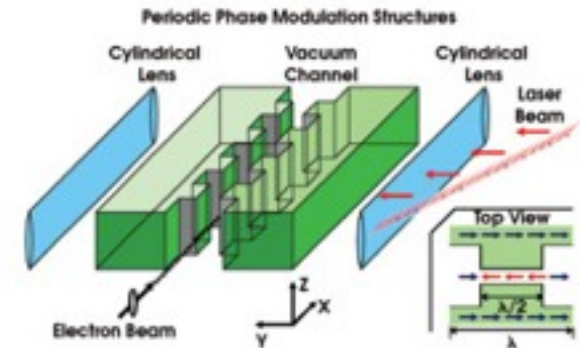
A laser-driven dielectric accelerator, i.e. photonic crystal accelerators, makes it possible to realize a tabletop/mobile system.

Structures of PCA

Phase Modulation Masked Structure

Silica, $\lambda=800\text{nm}$, $E_z=830\text{ MV/m}$

T. Plettner, *et al*, *PRST-AB*, **9**, 111301 (2006).



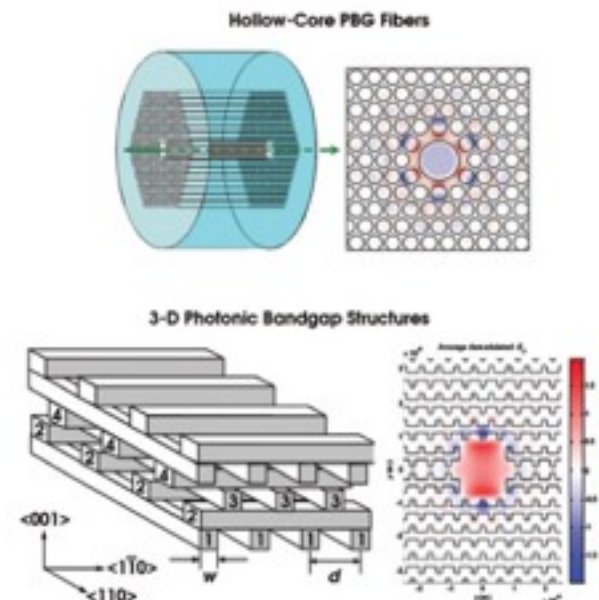
Wave Guiding Structures

Silica, $\lambda=1890\text{ nm}$, $E=400\text{ MV/m}$

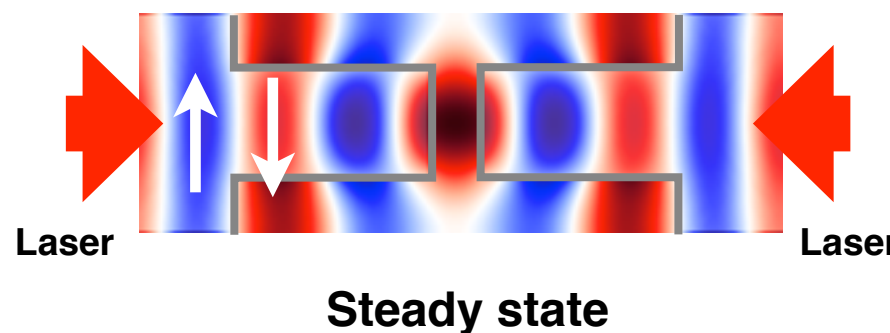
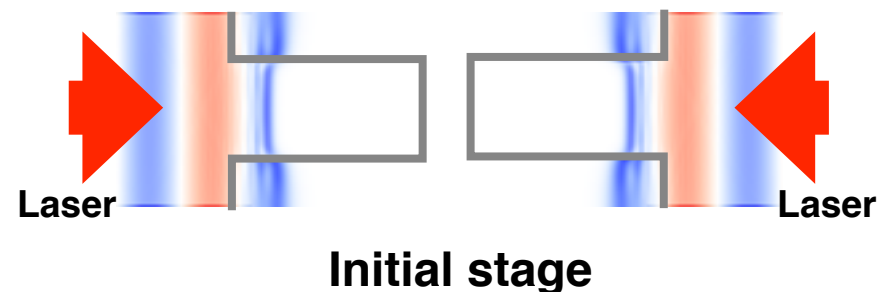
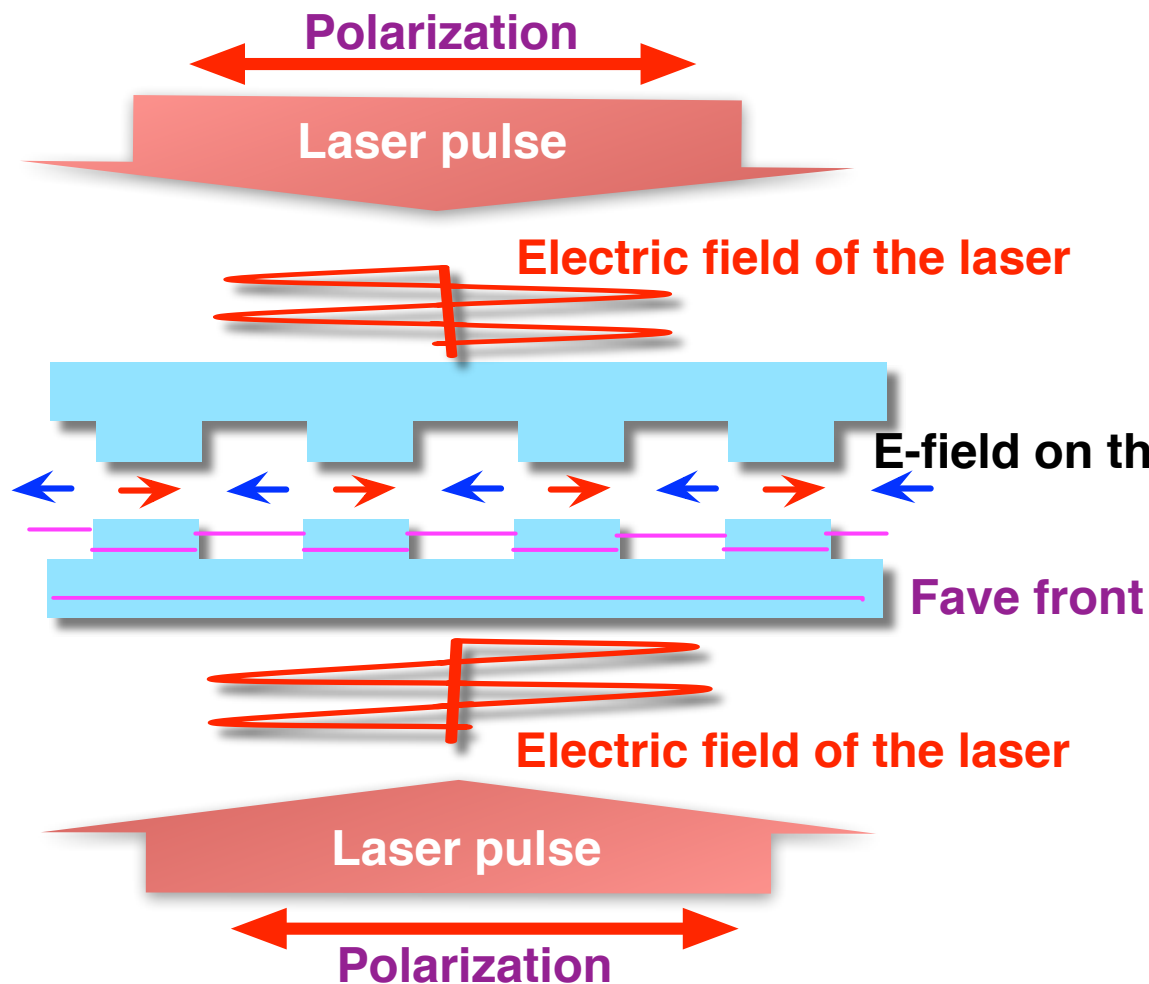
X. Lin, *PRST-AB*, **4**, 051301 (2001).

Silicon, $\lambda=2200\text{nm}$, $E_z=400\text{ MV/m}$

B. Cowan, *et al*, *PRST-AB*, **11**, 011301 (2008).



Alternate Direction of the Field is Produced by the Optical Path-difference



Summary of Analytical Consideration

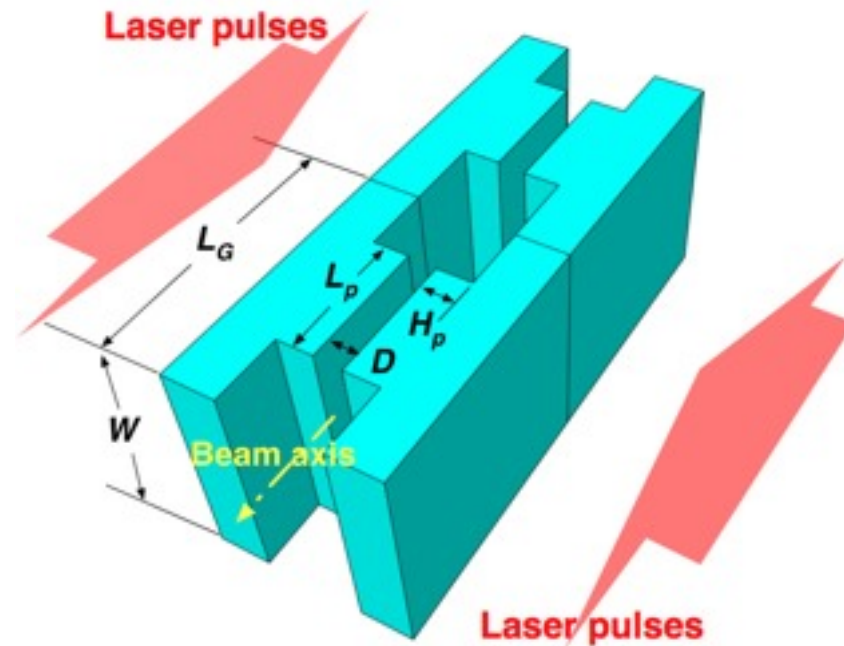
$$L_G = 2L_P$$

$$\frac{H_P}{\lambda_L} = \frac{1}{2(n-1)}$$

$$D \leq \frac{2L_P^2}{N\lambda_L}$$

$$L_G \leq \lambda_L$$

$$v_0/c = L_G/\lambda_L$$



If the initial electron is non-relativistic, the grating constant must be gradually changed from small value to the laser wavelength.

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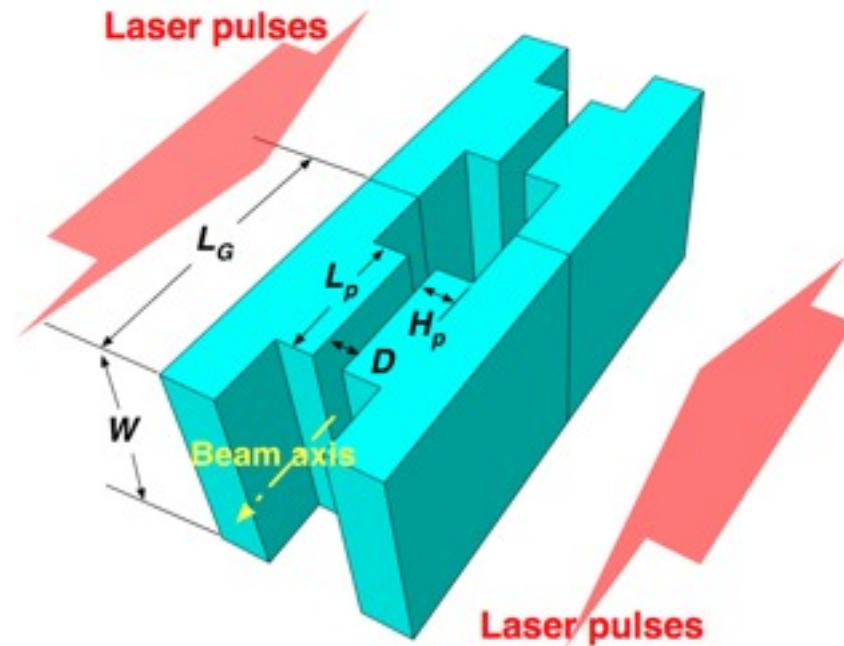
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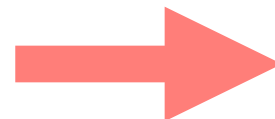
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In a real situation, the refractive and interference effects deform the field structure.

The asymmetry of the field, higher harmonics of the field relaxes the matching condition.

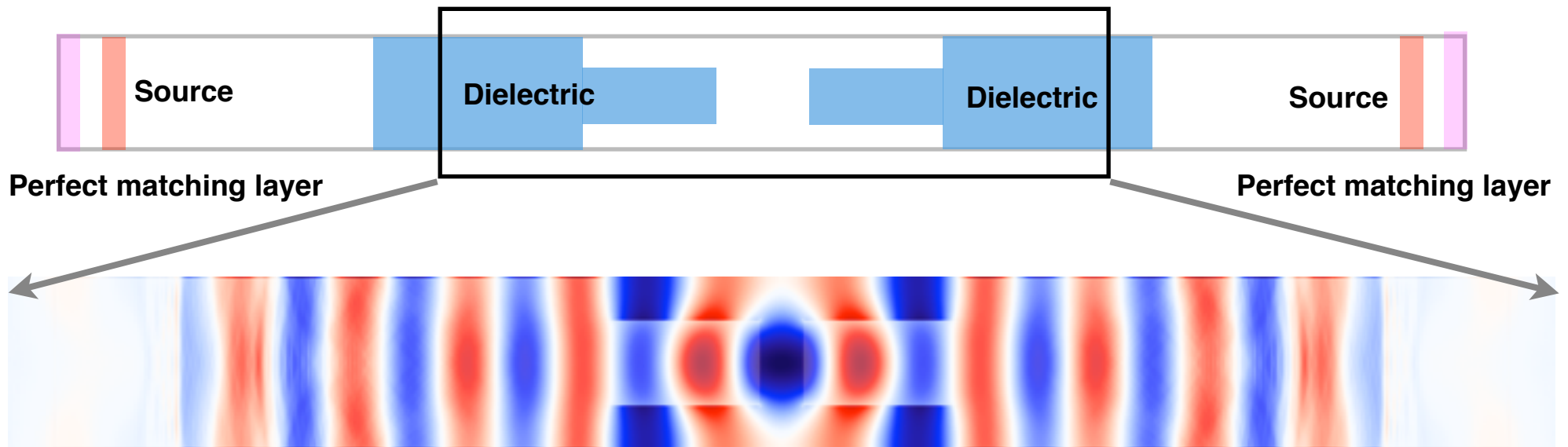


Numerical simulation

FDTD-simulation

FDTD (Finite-Difference Time-Domain method)
simulation software package developed at MIT
to model electromagnetic systems.

<http://ab-initio.mit.edu/wiki/index.php/Meep>



Intensity of the laser pulse ; $I_L = 10^{13} \text{ W/cm}^2$ (8.7GV/m) on entrance surfaces of dielectric.

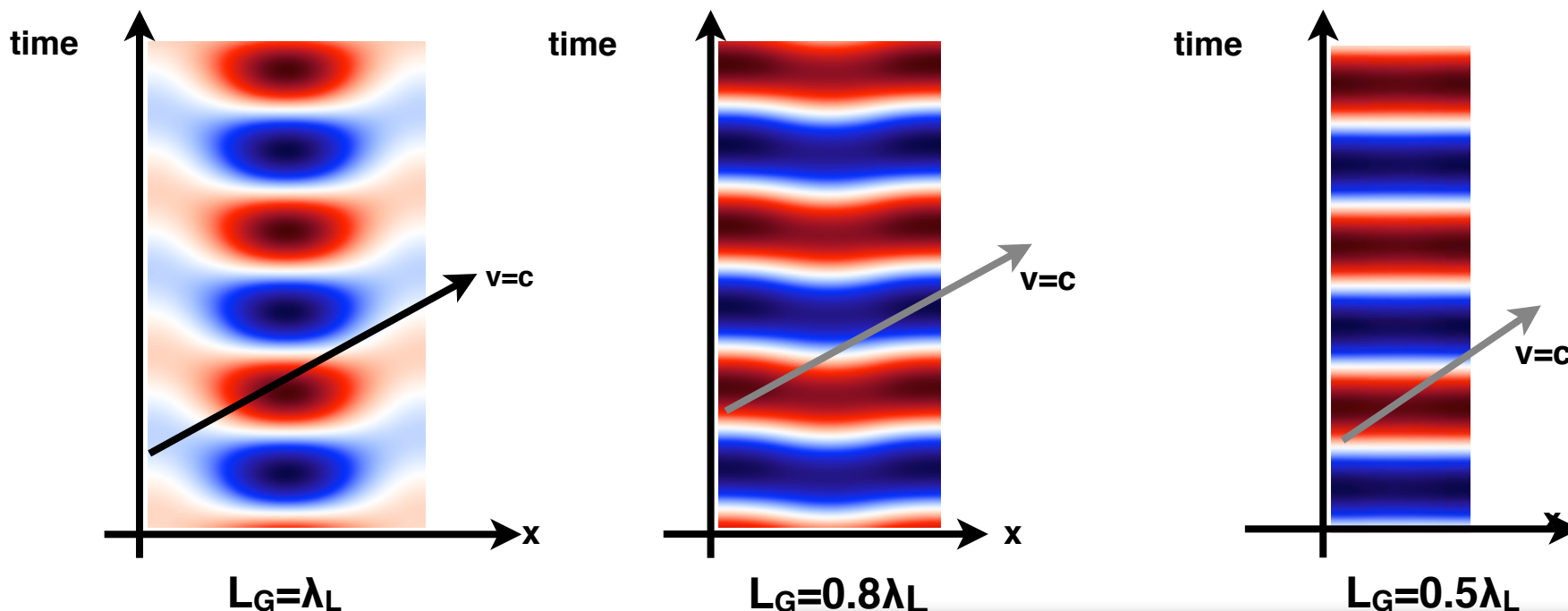
Laser wavelength ; $\lambda_L = 1.55 \mu\text{m}$.

Channel width ; $D = \lambda/4 = 0.39 \mu\text{m}$

Pillar height ; ($H_p = 0.9\lambda = 1.5 \mu\text{m}$) varied.

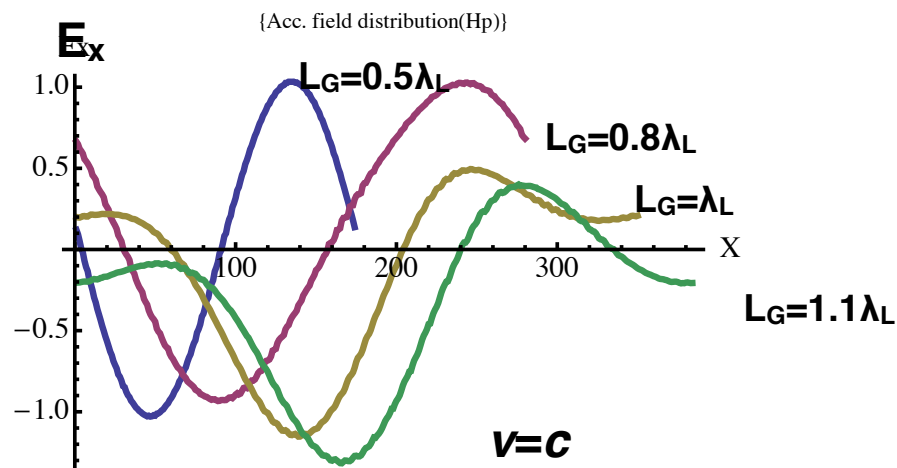
Acceleration field along the axis

Field strengths in xt -space for different grating constants



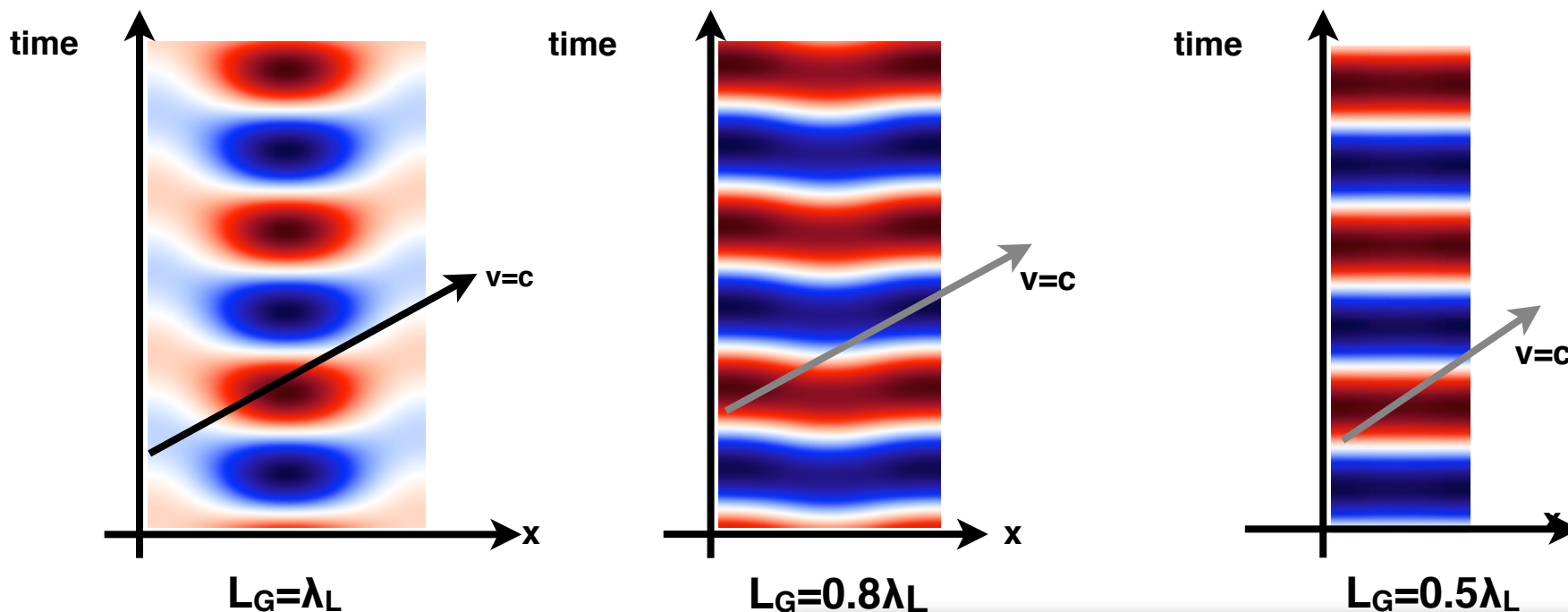
$$E_{acc} = \frac{1}{L} \int_0^L E_x(x, t) dx$$

$$x = vt + x_0$$



Acceleration field along the axis

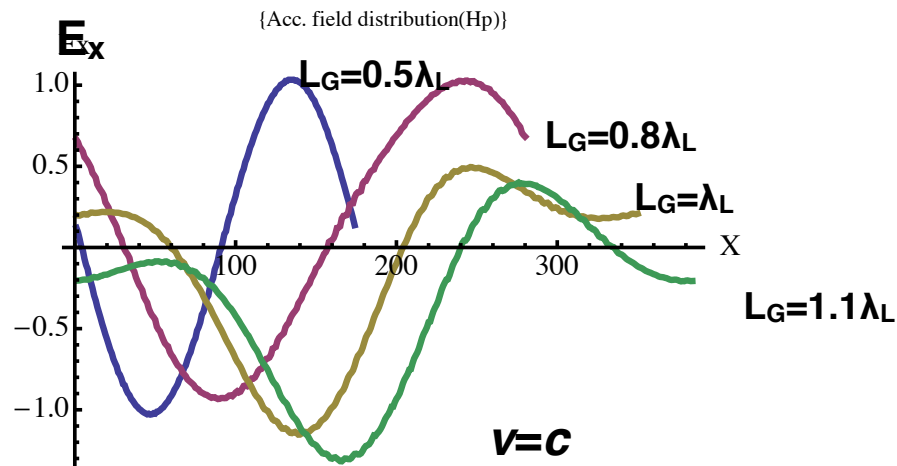
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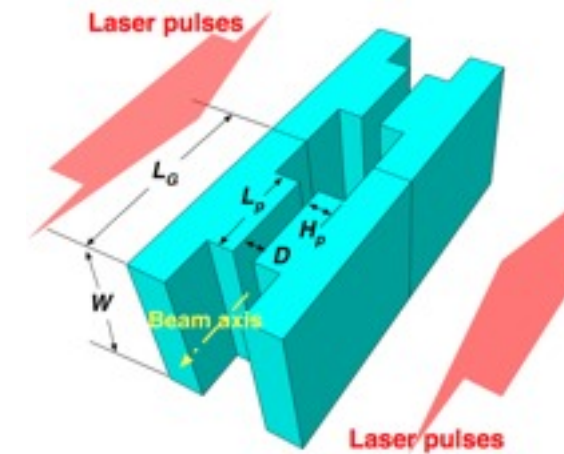
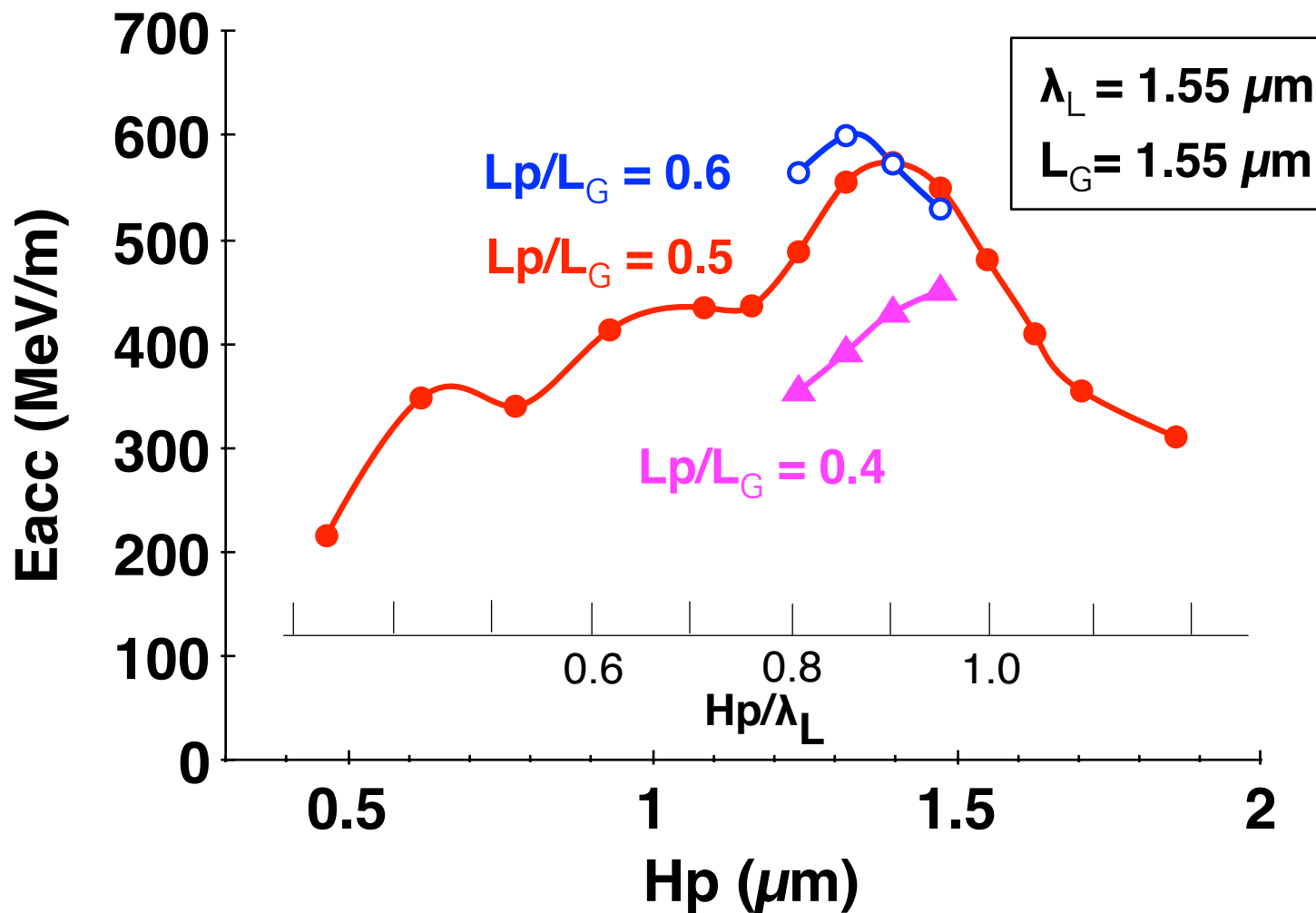
$$x = vt + x_0$$

very sensitive to the initial phase of the injection



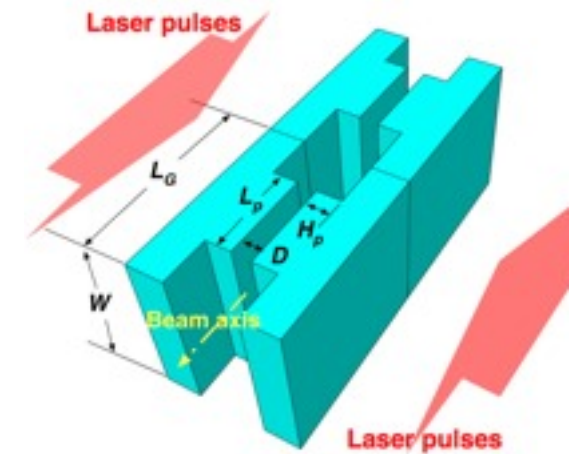
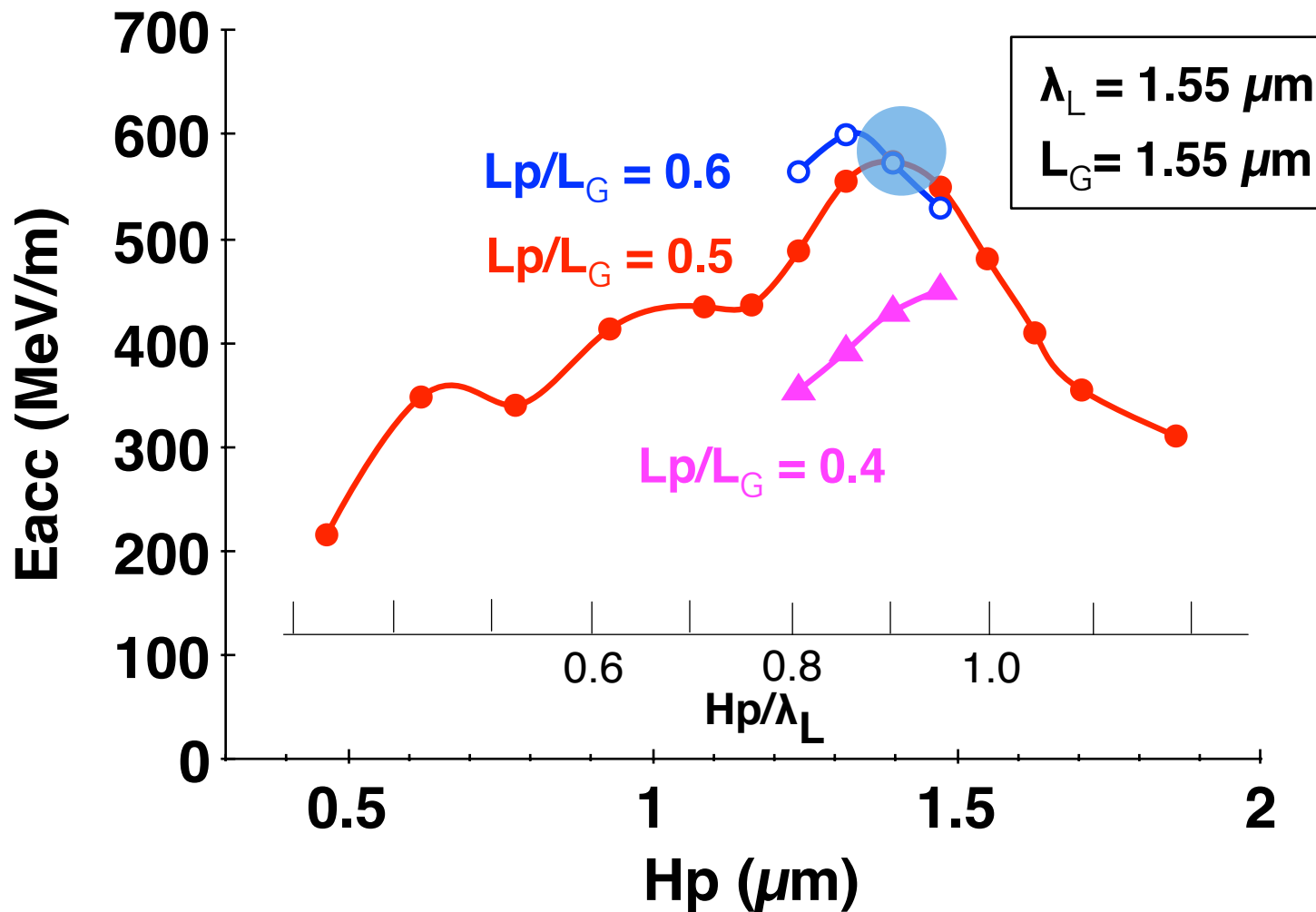
Optimum Pillar Height

$I_{th}=10^{13}$ W/cm² (1 J/cm @ 100fs) ; $E_{th}= 8.7$ GV/m is assumed.



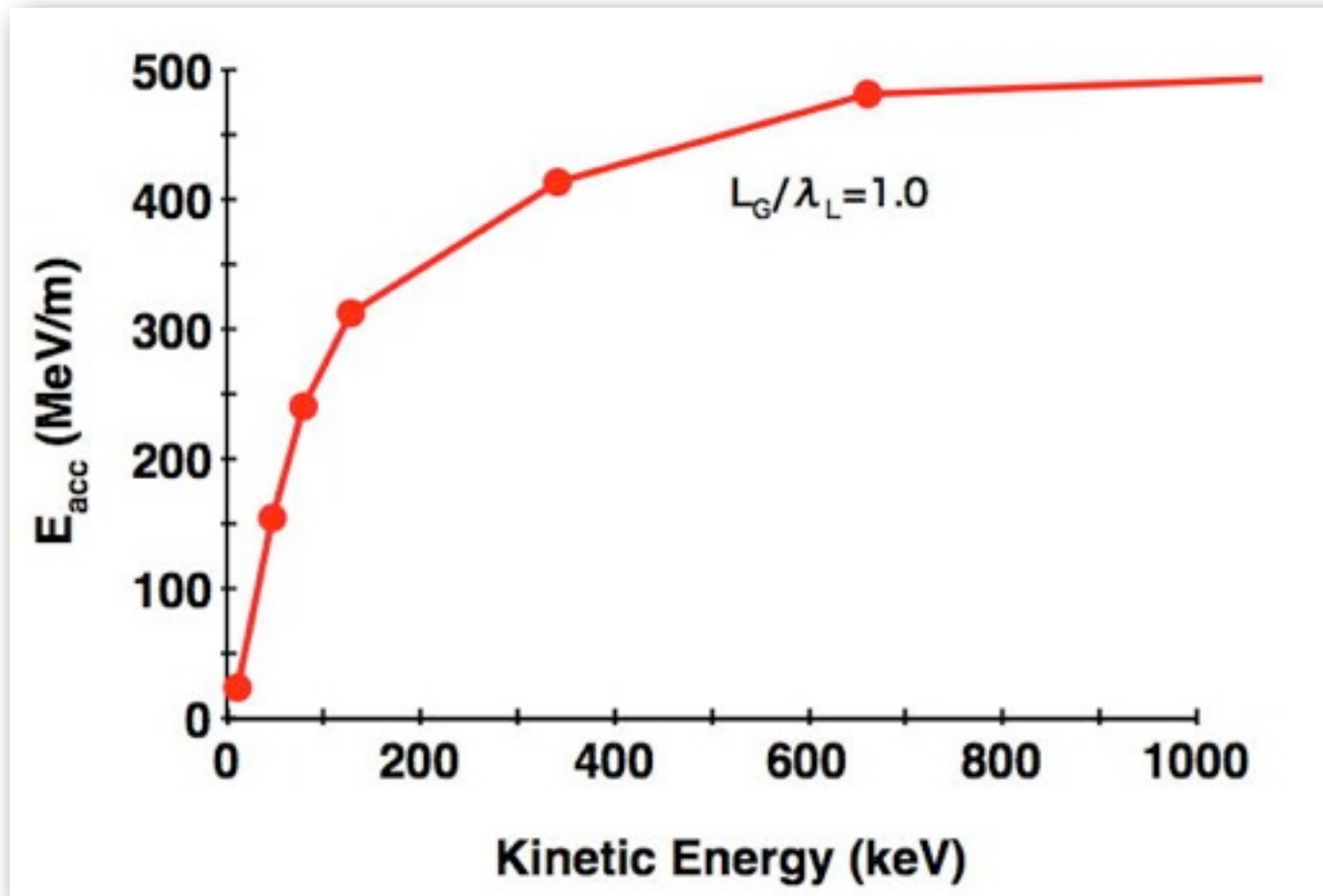
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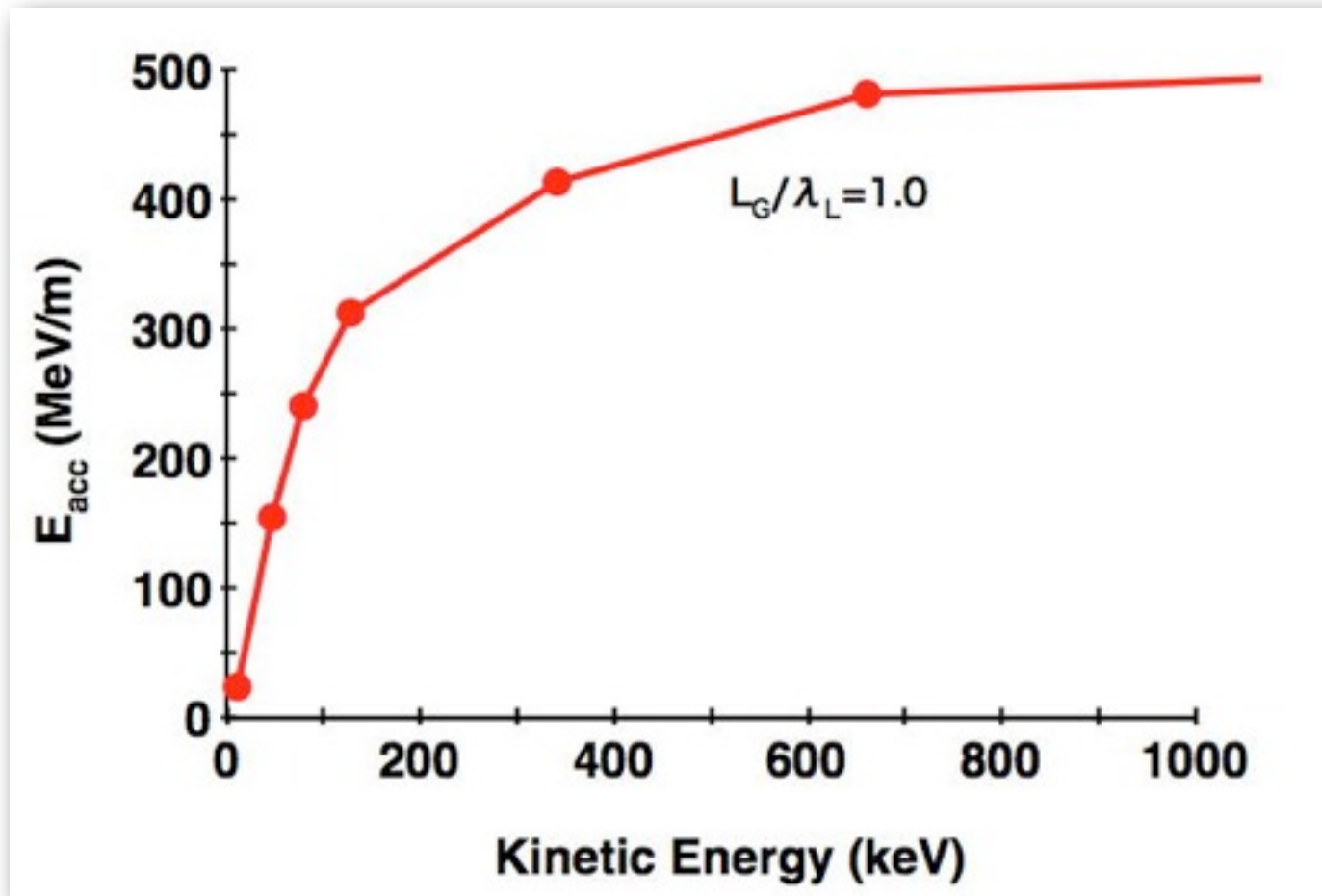


Optimum parameters of the structure are
 $L_p/L_G \approx 0.5$ and $H_p/\lambda_L \approx 0.9$.

Acceleration from the low energy in the constant grating period

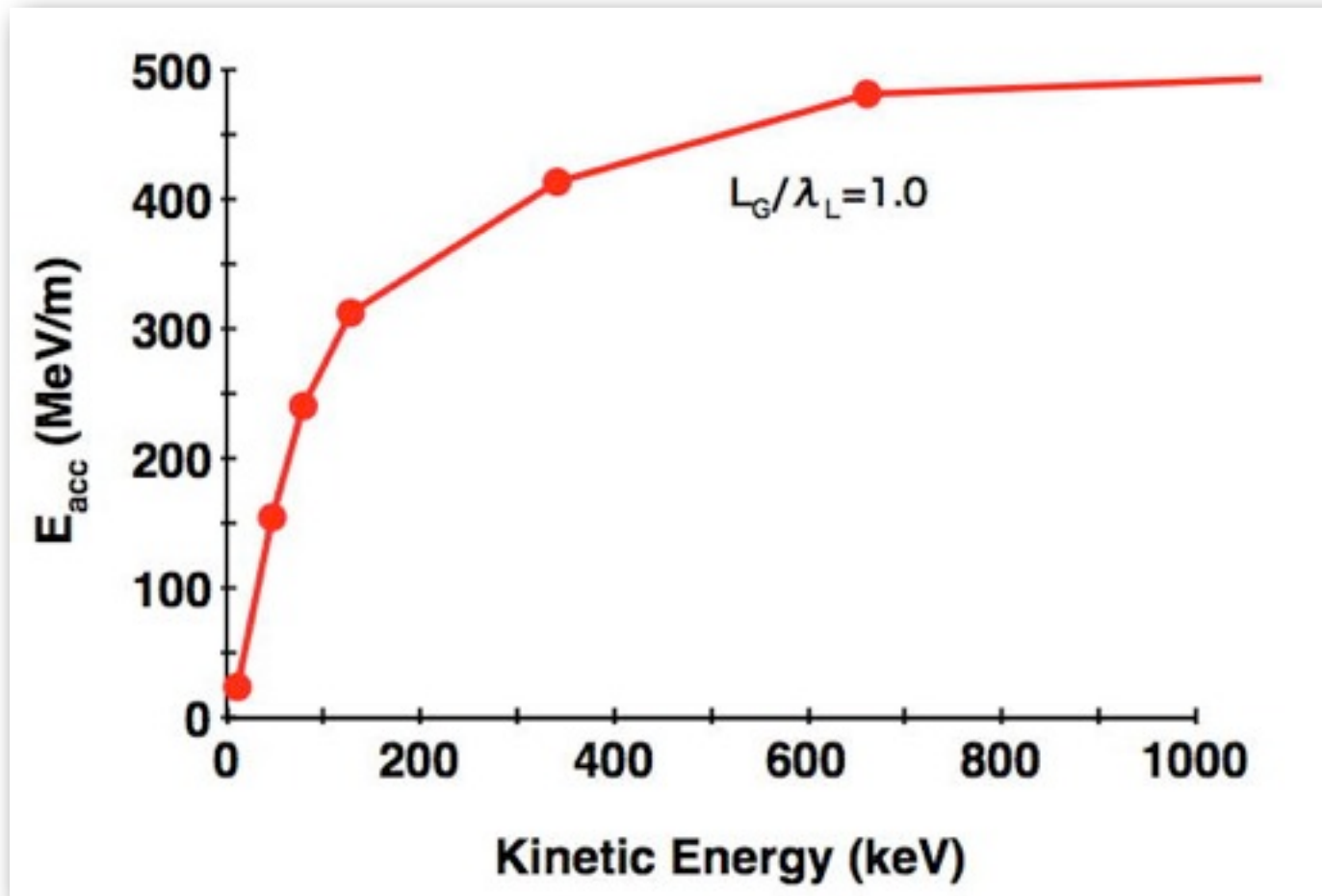


Acceleration from the low energy in the constant grating period



Slow electron can be accelerated even if the period is constant.

Acceleration from the low energy in the constant grating period

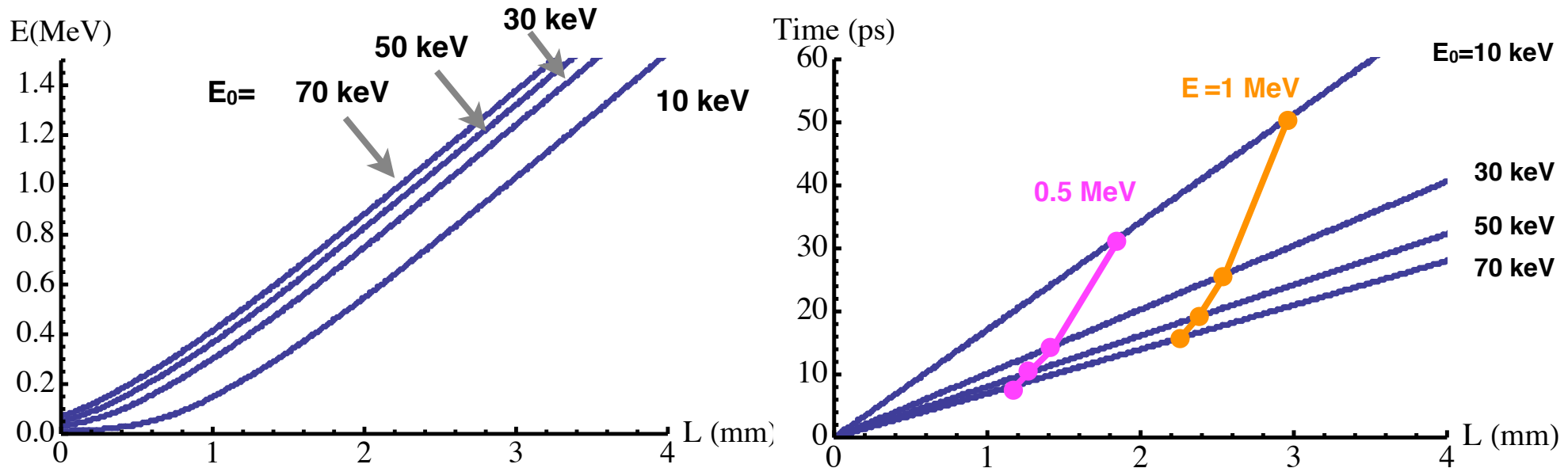


Slow electron can be accelerated even if the period is constant.

Acceleration field is very sensitive to the injection phase.

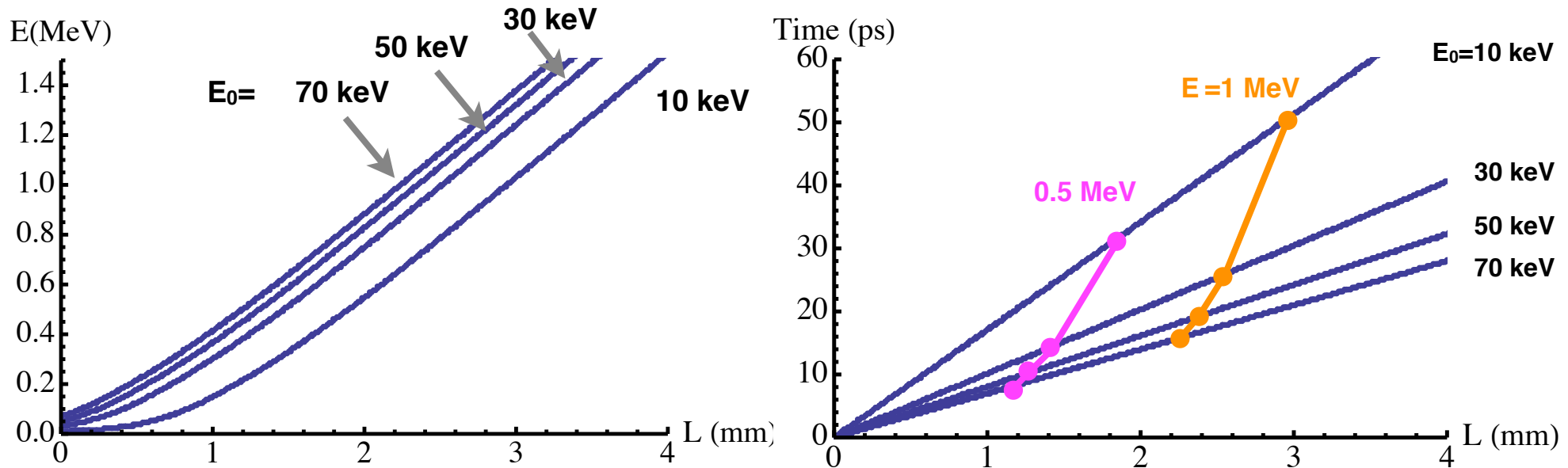
Acceleration Length and Time

Laser intensity = 10^{13} W/cm² (E = 8.7 GV/m)



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Accel. length to get 1 MeV,
3 mm for E₀ = 10 keV
2.2 mm for 70 keV

Accelerator and Laser Parameters

			1 MeV
Accelerator Length	L_A		3 mm
Channel width	W		0.1 mm
Irradiation Area (per side)	$A = W L_A$		$3 \times 10^{-3} \text{cm}^2$
Laser Intensity (damage threshold)	I_{th}		10^{13}W/cm^2
Two sides irradiation	Laser Power (total)	$P_L = 2P_{th}A$	60 GW
	Pulse width	$\tau_L = L_A/v_{eff}$	50 ps
	Energy (total energy)	$E_L = P_L\tau_L$	3 J

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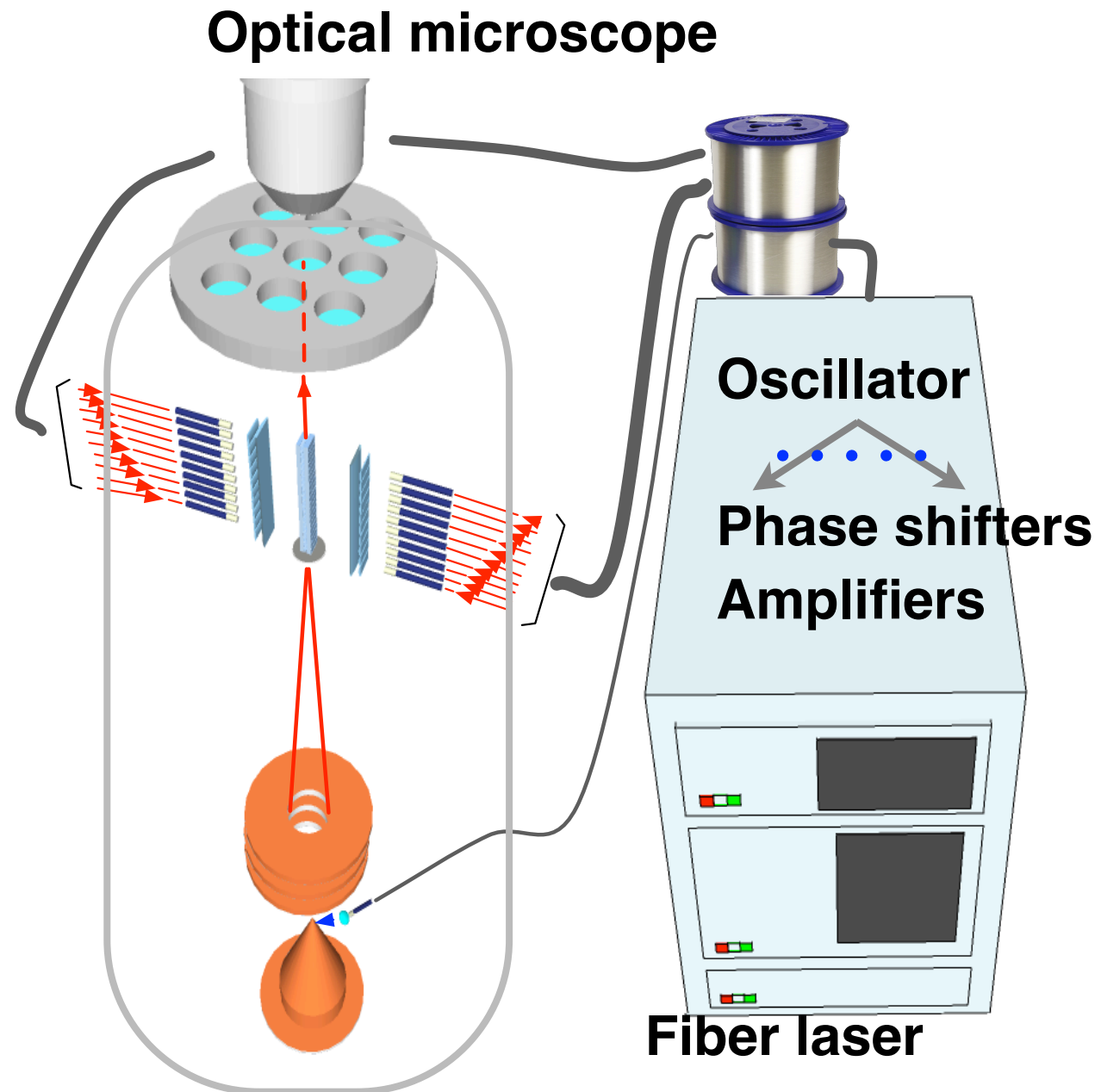
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Accelerator Length	L_A		3 mm
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In order to decrease the laser energy, the laser pulse must locally irradiate around the electron bunch.

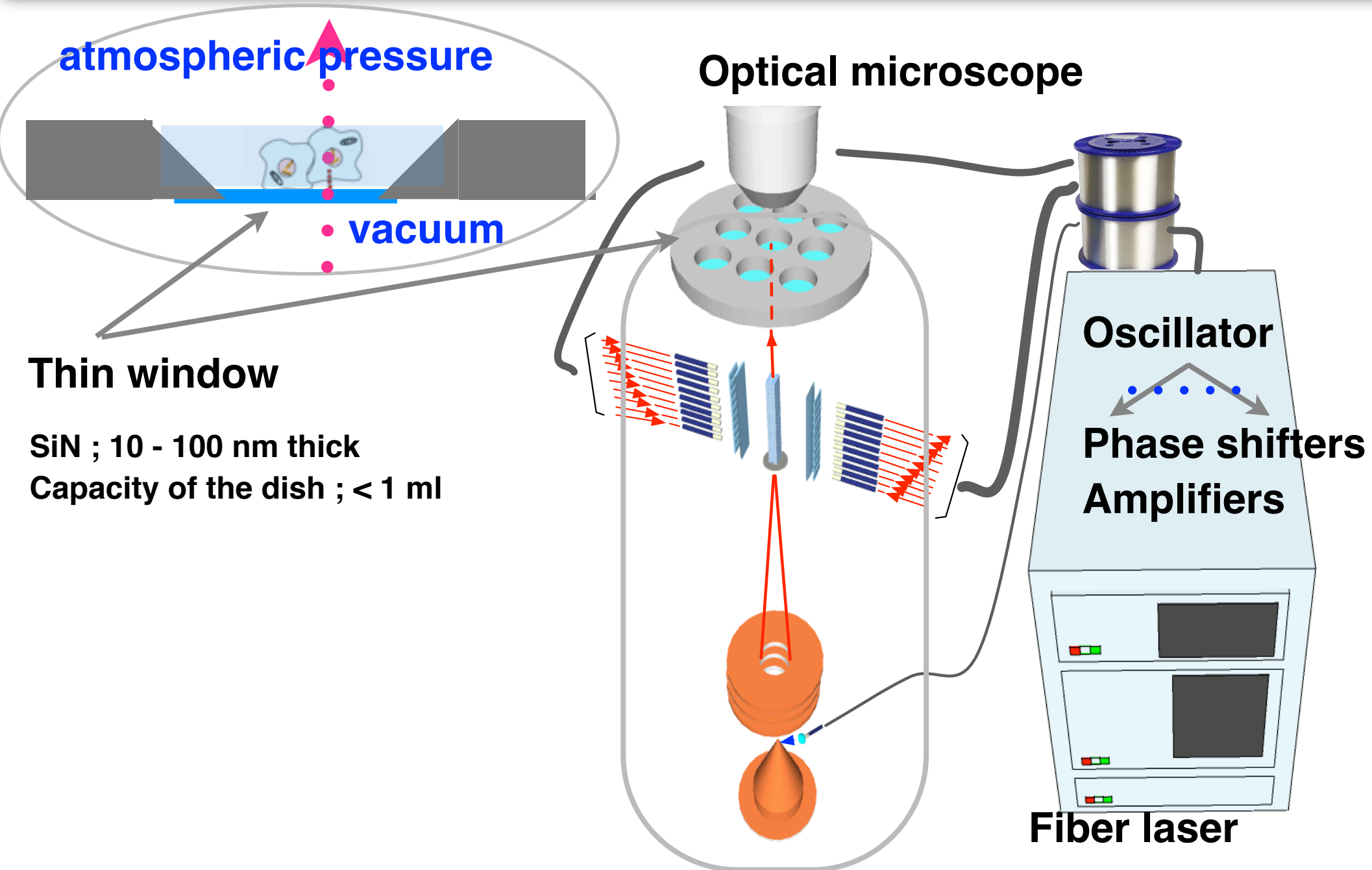
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	Pulse width	$\tau_L = L_A/v_{eff}$	50 ps
	Energy (total energy)	$E_L = P_L\tau_L$	3 J
Divided irradiation	Number of Pulses (one side)	N	10
	Laser Power per Pulse	$P_{L,N} = P_L/N$	3 GW
	Pulse Width	$\tau_{L,N} = \tau_L/N$	5 ps
	Energy (total)	$E_{L,N} = E_L/N$	30 mJ
	(per pulse)	$E_{p,N} = E_L/2N^2$	1.5 mJ

Sketch of a fiber-laser-pumped accelerator



Sketch of a fiber-laser-pumped accelerator



Summary

1. The deformation of the wavefront in the phase-modulation-masked-type accelerator relaxed the matching condition. The **slow electron** to be accelerated even under the geometry of $L_G = \lambda_L$.
2. The electron initially at the low energy of **20 keV** felt the acceleration field strength of **20 MV/m** and gradually felt higher field as the speed was increased. The ultra **relativistic** electron felt the field strength of **600 MV/m**.
3. Restrictions on the laser is relaxed by adopting sequential laser pulses. The required laser power is estimated to be **3 GW/pulse** when ten pairs of sequential laser pulse is irradiated.